

# COMPUTER PROGRAM SPM-MC AND ITS APPLICATIONS IN EMR STUDIES OF TRANSITION IONS IN CRYSTALS

IRENEUSZ STEFANIUK<sup>1\*</sup>, CZESŁAW RUDOWICZ<sup>2</sup>

<sup>1</sup> Institute of Physics, University of Rzeszów, Rejtana 16a, 35-310 Rzeszów, Poland

<sup>2</sup> Institute of Physics, West Pomeranian University of Technology, Al. Piastów 17, 70-310 Szczecin, Poland

Received June 20, 2010; accepted June 29, 2010; published online November 20, 2010.

**Development of computer program for superposition model (SPM) calculations is reported. A novel aspect in this approach is the usage of the Monte Carlo (MC) method to predict the feasible ligands' positions by fitting the experimental electron magnetic resonance (EMR) data to theoretical estimates. The program SPM-MC may be helpful for interpretation of EMR spectra for paramagnetic ions at low (triclinic) symmetry sites and coordinated by up to 13 ligands of arbitrary type. For illustration we consider the case of chromium Cr<sup>3+</sup> ions in YAlO<sub>3</sub> (YAP) crystal, which exhibits orthorhombic structure (*Pbnm* space group), whereas the site symmetry is lower than orthorhombic due to structural distortions.**

## INTRODUCTION

Interpretation of electron magnetic resonance (EMR) spectra is often cumbersome due to low local site symmetry exhibited by paramagnetic centers. For such cases comparative analysis of data obtained from experimental EMR studies and the structural data may be facilitated by theoretical modeling of the zero-field splitting (ZFS) parameters (ZFSPs) - for a review, see, e.g. (Rudowicz, 1987; Rudowicz & Misra, 2001). Various modeling techniques for analysis and interpretation of EMR data for transition ions at low symmetry sites in crystals have recently been discussed in (Rudowicz & Gnutek, 2009). One major technique is the superposition model (SPM) (Newman & Urban, 1975; Newman & Ng, 1989), which enables prediction of ZFSPs based on the assumed structural distortions models for a transition ion (M) surrounded by the nearest neighbour *n* ligands (L) within ML<sub>*n*</sub> complex, see, e.g. (Rudowicz, 1987; Yeom, Chang, Choh & Rudowicz, 1994). Comparison of SPM predictions with experimental ZFSP values enables verification of the structural models and thus provides information about the local distortions around paramagnetic ions.

In this paper we present a computer program SPM-MC for modeling and analysis of the structure of ML<sub>*n*</sub> complexes around paramagnetic centers, which utilizes as input the experimental data extracted from EMR spectra. This program is based on the SPM technique and may also be used for prediction of the ZFSPs based on the structural data. A novel aspect in this approach is the usage of the Monte Carlo (MC) method to predict the feasible ligands positions by fitting the experimental EMR data to theoretical estimates. The program has been tested for various types of crystal structures and ML<sub>*n*</sub> complexes.

Subsequently, we consider its application for SPM analysis for Cr<sup>3+</sup> and Fe<sup>3+</sup> ions in yttrium aluminum perovskite YAlO<sub>3</sub> (YAP) crystal. The EMR measurements of YAP single crystals containing Cr<sup>3+</sup> and Fe<sup>3+</sup> impurity ions as well as the spin Hamiltonian analysis of EMR spectra (carried out based on triclinic C<sub>*i*</sub>) site symmetry were reported in (Stefaniuk, Rudowicz, Gnutek & Suchocki, 2009). Capabilities of the program SPM-MC and preliminary results of its application for Cr<sup>3+</sup> ions in YAP are presented here.

## SPIN HAMILTONIAN AND SPM ANALYSIS OF EMR SPECTRA

In this study we adopt the spin Hamiltonian (SH) notation (Rudowicz, 1987; Rudowicz & Misra, 2001), outlined in (Stefaniuk, Rudowicz, Gnutek & Suchocki, 2009). For clarity we only define the form of SH suitable for arbitrary low (triclinic) symmetry and expressed in terms of the extended Stevens operators (ESO)  $O_k^q$  defined in (Rudowicz, 1985; Rudowicz & Chung, 2004):

$$H_s = \mu_B \mathbf{B} \cdot \mathbf{g} \cdot \mathbf{S} + \sum B_k^q O_k^q(S_x, S_y, S_z) \\ = \mu_B \mathbf{B} \cdot \mathbf{g} \cdot \mathbf{S} + \sum f_k b_k^q O_k^q \quad (1)$$

The ZFSPs  $B_k^q$  ( $b_k^q$ ) with  $k = 2$  only are required for Cr<sup>3+</sup> ( $S=3/2$ ), whereas those with  $k = 2$  and 4 for Fe<sup>3+</sup> ( $S=5/2$ ) ions (Rudowicz, 1987; Rudowicz & Misra, 2001; Stefaniuk, Rudowicz, Gnutek & Suchocki, 2009). For proper relations between the orthorhombic ZFSPs and the conventional ones (Rudowicz, 1987; Rudowicz & Misra, 2001) as well as pertinent

\*Corresponding author: <istef@univ.rzeszow.pl>

references, the readers may refer to (Stefaniuk, Rudowicz, Gnutek & Suchocki, 2009; Rudowicz, 2000).

Low symmetry aspects observed in EMR spectra of  $\text{Cr}^{3+}$  and  $\text{Fe}^{3+}$  ions in YAP arise due to structural distortions. YAP has a distorted perovskite structure with the orthorhombic centrosymmetric  $Pbnm$  space group (Geller & Wood, 1956; Vasylechko L., Matkovskii A., Savytski D., Suchocki A. & Wallrafen F., 1999). The YAP structure forms as a grid of tilted  $\text{AlO}_6$  regular octahedra with Y ions occupying voids between them. The  $\text{YO}_8$  dodecahedra are considerably distorted since the Y-O bond-length varies from 0.2284 nm to 0.2597 nm for YAP-Nd (Vasylechko *et al.*, 1999). The cation ionic radii are:  $R_Y = 0.102$  nm and  $R_{\text{Al}} = 0.053$  nm. The  $\text{Cr}^{3+}$  ( $R_{\text{Cr}} = 0.062$  nm),  $\text{Ti}^{3+}$  ( $R_{\text{Ti}} = 0.067$  nm),  $\text{Mn}^{4+}$  ( $R_{\text{Mn}} = 0.053$  nm), and  $\text{Fe}^{3+}$  ( $R_{\text{Fe}} = 0.065$  nm) ions are found to substitute the Al cations in YAP host (Yamaga M., Takeuchi H., Han T. J. & Henderson B., 1993; Belt *et al.*, 1974; Rakhimov R. R., Wilkerson A. L., Loutts G. B., Noginov M. A., Noginova N., Lindsay W. & Ries H. R., 1998; Yamaga M., Yosida T., Henderson B., O'Donnell K. & Date M., 1992).

Superposition model represents the ZFSPs in the form of linear combinations of products of the intrinsic parameters and coordination factors – for definitions, see, e.g. (Newman & Urban, 1975; Newman & Ng, 1989; Rudowicz, 1987; Yeom, Chang, Choh & Rudowicz, 1994). The intrinsic parameters depend only on the kind of ligands and their distances from the central ion. The coordination factors  $K_k^q(\theta, \phi_i)$  depend only on the angular positions of ligands in the coordinate system associated with the paramagnetic ion. Full listing of  $K_k^q(\theta, \phi_i)$  obtained using the transformation properties of the ESOs for arbitrary symmetry and  $k = 2, 4$ , and  $6$  has been provided in (Rudowicz, 1987). For illustration we provide the SPM expressions used for the orthorhombic ZFSPs  $b_k^q$  suitable for  $\text{Cr}^{3+}$  ions (Müller & Berlinger, 1983; Müller, Berlinger & Albers, 1985; Kuriata J., Baker J. M., Sadlowski L., Stefaniuk I. & Bodziony T., 1998):

$$b_2^0 = \sum_i \left( -A \left( \frac{R_0}{R_i} \right)^n + B \left( \frac{R_0}{R_i} \right)^m \right) \cdot \frac{1}{2} (3 \cos^2 \theta - 1) \quad (2)$$

$$b_2^2 = \sum_i \left( -A \left( \frac{R_0}{R_i} \right)^n + B \left( \frac{R_0}{R_i} \right)^m \right) \cdot \frac{3}{2} (\sin^2 \theta \cos 2\phi)$$

and  $\text{Fe}^{3+}$  ions (Siegel & Müller, 1979; Müller & Berlinger, 1983):

$$b_2^0 = \sum_i \left( (-A + B) \left( \frac{R_0}{R_i} \right)^{t_2} \right) \cdot \frac{1}{2} (3 \cos^2 \theta - 1) \quad (3)$$

$$b_2^2 = \sum_i \left( (-A + B) \left( \frac{R_0}{R_i} \right)^{t_2} \right) \cdot \frac{3}{2} (\sin^2 \theta \cos 2\phi)$$

The program (Section 3.1) handles also the general SPM expressions for  $k = 2$  and  $4$  (Newman & Ng, 1989; Rudowicz, 1987; Gnutek, Acikgoz & Rudowicz, 2010):

$$b_k^q = \sum_i \bar{b}_k(R_i) \cdot K_k^q(\theta, \phi_i),$$

$$\bar{b}_k(R_i) = \bar{b}_k(R_0) \left( \frac{R_0}{R_i} \right)^{t_k} \quad (4)$$

In Eqs. (2)-(4) ( $\theta, \phi$ ) are the polar angles,  $R_i$  - the distance between the central ion and  $i^{\text{th}}$ -ligand; the constants  $A, B, n, m, t_k, \bar{b}_k$  and the reference distance  $R_0$  depend on the kind of impurity ions and their valence (Rudowicz, 1987; Müller & Berlinger, 1983; Müller, Berlinger & Albers, 1985; Kuriata *et al.*, 1998; Siegel & Müller, 1979). The SPM calculations for the 2<sup>nd</sup>-rank ZFS terms may be carried out in two ways: (i) based on the method (Müller & Berlinger, 1983) using Eqs. (2) or (3) and (ii) using general expressions, Eqs. (4), for all 5 triclinic ZFSPs. For the 4<sup>th</sup>-rank ZFS terms Eqs. (4) are used together with the selected way for the 2<sup>nd</sup>-rank ZFS terms. Development of the extended version of the program SPM-MC based on the general SPM expressions for  $k = 6$ , see, e.g. (Gnutek, Acikgoz & Rudowicz, 2010) and references therein, is in progress.

## PROGRAM SPM-MC AND ITS APPLICATIONS

### Program SPM-MC

Basic capabilities of the program SPM-MC are outlined below. The input data, i.e. the unit cell parameters and the number of ligands is first entered in the field 'Structural data' (see, Fig. 1). Next, in the field 'Ligands' coordinates' the positions of up to 13 distinct ligands may be entered. The experimental EMR data are provided in the field 'ZFS parameters', whereas the rank of ZFSPs to be used for calculations is selected via 'Rank considered'. The values of the model parameters ( $A, B, n, m, t_k, \bar{b}_k, R_0$ ) are entered in the field 'Model parameters'. At present two calculation options, i.e. 'Version' #1 and #2, are available based on the relations (2) or (3) and (4), respectively. The required range and accuracy of calculations are specified in the field 'Computational options'. Finally, the program may be activated by clicking on the button 'Start MC'. The results of calculations based on the Monte Carlo (MC) method represented in the spherical and Cartesian coordinates, together with the values of all EMR parameters, are

displayed in the two upper-right windows. Additionally, the graphical 3D representation of the ligands positions and the axis systems appears in the lower-right window. Program enables to obtain one

set of the values of the ZFSPs  $b_k^q$  by clicking on the button 'Run for one set'. All numerical output results may be saved into a text file, whereas the graphs are saved as \*.bmp files.

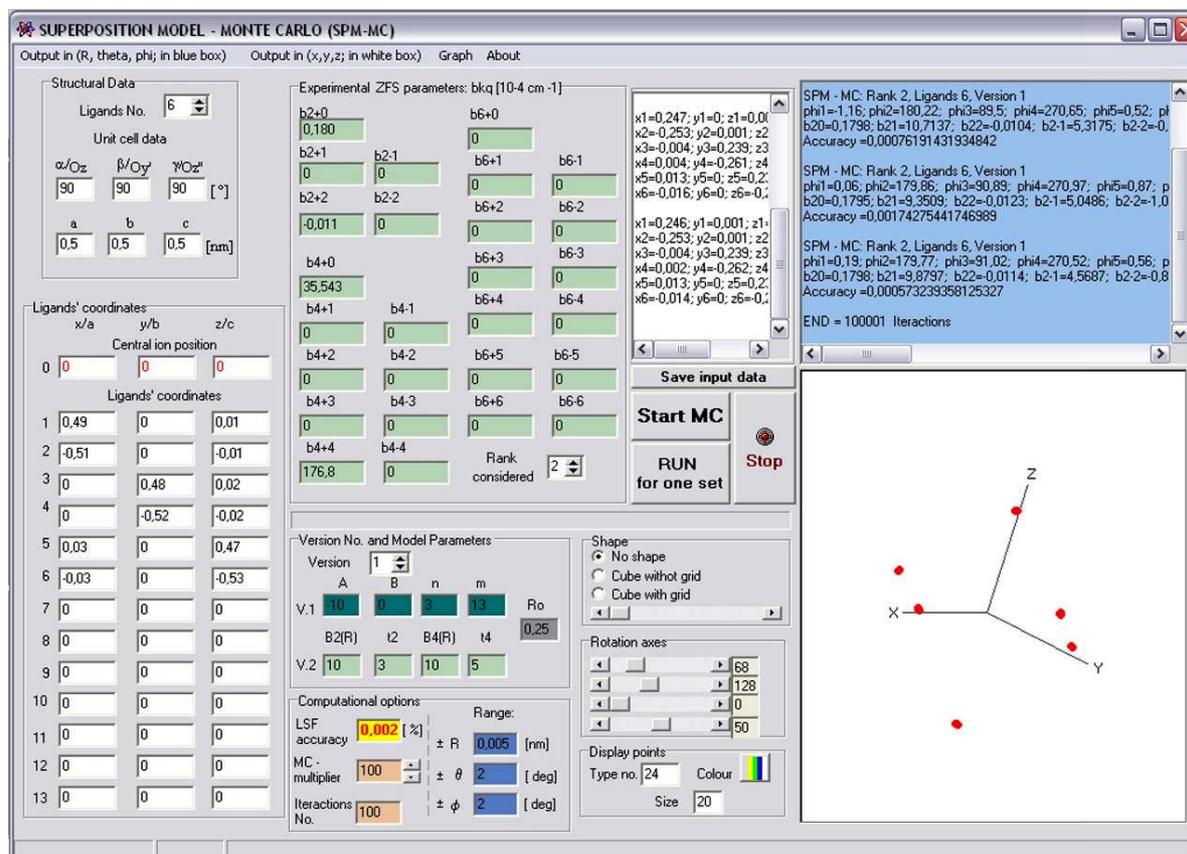


Fig. 1. Sample interface for the program SPM-MC showing illustrative input and output data.

### Applications

For  $\text{Cr}^{3+}$  ions, which replace  $\text{Al}^{3+}$  in  $\text{YAlO}_3$ , the parameters  $A$ ,  $B$ , and  $R_0$  were obtained from the data in (Müller & Berlinger, 1983; Siegel & Müller, 1979), whereas the crystallographic data from (Diehl & Brandt, 1975). For illustration, we adopt the following values:  $R_0 = 0.195$  nm,  $n = 10$ ,  $m = 13$ ,  $A = -10.6$   $\text{cm}^{-1}$ , and  $B = -8.2$   $\text{cm}^{-1}$ . Using the ZFS parameters and the pertinent conversion relations

provided in (Stefaniuk, Rudowicz, Gnutek & Suchocki, 2009; Rudowicz, 2000) the program SPM-MC computes the feasible positions of the oxygen ligands in a given unit cell volume that yield the SPM-predicted ZFSPs consistent with the experimental ZFSPs. The numerical calculations and analysis were carried out for the complex  $\text{Cr}(\text{Al})\text{O}_6$  in  $\text{YAlO}_3$ .

Table 1. The positions of the oxygen ligands ( $i = 1$  to 6) in the  $\text{MO}_6$  octahedra in the polar coordinates.

Ligand's number $i$ :		1	2	3	4	5	6
Host $\text{AlO}_6$ (Diehl & Brandt, 1975)	$R_i$ [nm]	0.1901	0.1901	0.1910	0.1910	0.1921	0.1921
	$\theta_i$ [°]	14.2	165.8	80.2	99.8	80.2	99.8
	$\phi_i$ [°]	196.7	16.7	56.0	236.0	324.4	144.4
$\text{Cr}(\text{Al})\text{O}_6$ our work	$R_i$ [nm]	0.2010	0.1967	0.2041	0.2033	0.2088	0.2162
	$\theta_i$ [°]	16.3	86.5	88.1	101.2	175.6	100.1
	$\phi_i$ [°]	205.1	57.2	331.7	241.6	22.0	154.1

The angles  $\theta_i$  and  $\phi_i$  were varied in the range of several degrees from the starting values reported for the undistorted  $\text{AlO}_6$  octahedron, whereas the minimum distances  $R_i$  were computed from the ionic radii. The averaged values ( $\theta_i$ ,  $\phi_i$ ,  $R_i$ ) obtained from the program SPM-MC are listed in Table 1 together with the crystallographic data for the  $\text{AlO}_6$  octahedron in pure crystal. These values were used to construct the  $\text{Cr}(\text{Al})\text{O}_6$  octahedron in  $\text{YAlO}_3$  (not shown here).

## CONCLUSIONS

Computer program SPM-MC based on the superposition model (SPM) approach and incorporating the Monte Carlo (MC) method has been developed for interpretation of EMR spectra for paramagnetic ions coordinated by up to 13 ligands of arbitrary type. The extended version of this program will be suitable for low (triclinic) symmetry sites. This program enables prediction of the feasible ligands' positions by fitting the experimental EMR data to the theoretical estimates. Preliminary results for the impurity  $\text{Cr}^{3+}$  ions in  $\text{YAlO}_3$  are presented. SPM analysis confirms the most probable model of distortions around  $\text{Cr}^{3+}$  ions occupying the Al positions. The predicted small radial distortions of the  $\text{Cr}(\text{Al})\text{O}_6$  complex correlate well with the size of the ionic radii of the dopant ions.

## ACKNOWLEDGEMENTS

This work was partially supported by a grant from the Polish Ministry of Science and Education for the years 2009-2010 to CZR. Thanks are due to Mr P. Gnutek for helpful comments.

## REFERENCES

- Belt R. F., Latore J. R., Uhrin R. & Paxton J. (1974). EPR and optical study of Fe in Nd:YAlO<sub>3</sub> laser crystals. *Appl. Phys. Lett.* **25**, 218-220.
- Diehl R. & Brandt G. (1975). Crystal structure refinement of YAlO<sub>3</sub>, a promising laser material. *Mater. Res. Bull.* **10**, 85-90.
- Geller S. & Wood E. A. (1956). Crystallographic studies of perovskite-like compounds. I. Rare earth orthoferrites and YFeO<sub>3</sub>, YCrO<sub>3</sub>, YAlO<sub>3</sub>. *Acta Crystallogr.* **9**, 563-568.
- Gnutek P., Acikgoz M. & Rudowicz C. (2010). Superposition model analysis of the zero-field splitting parameters of Fe<sup>3+</sup> doped in TlInS<sub>2</sub> crystal – Low symmetry aspects. *Opt. Mat.*, **32**, 1161-1169.
- Kuriata J., Baker J. M., Sadlowski L., Stefaniuk I. & Bodziony T. (1998). Superposition model analysis of Cr<sup>3+</sup> in sodium ammonium sulphate dehydrate. *J. Phys.: Condens. Matter* **10**, 407-415.
- Müller K. A. & Berlinger W. (1983). Superposition model for sixfold-coordinated Cr<sup>3+</sup> in oxide crystals (EPR study). *J. Phys. C: Solid State Phys.* **16**, 6861-6874.
- Müller K. A., Berlinger W. & Albers J. (1985). Paramagnetic resonance and local position of Cr<sup>3+</sup> in ferroelectric BaTiO<sub>3</sub>. *Phys. Rev. B* **32**, 5837-5844.
- Newman D. J. & Ng B. (1989). The superposition model of crystal fields *Rep. Prog. Phys.* **52**, 699-763.
- Newman D. J. & Urban W. (1975). Interpretation of S-state ion E.P.R. spectra. *Adv. Phys.* **24**, 793-844.
- Rakhimov R. R., Wilkerson A. L., Loutts G. B., Noginov M. A., Noginova N., Lindsay W. & Ries H. R. (1998). Spin and valence states of manganese ions in manganese-doped yttrium orthoaluminate. *Solid State Communications*, **108**, 8 549-554.
- Rudowicz C. & Chung C. Y. (2004). The generalization of the extended Stevens operators to higher ranks and spins, and a systematic review of the tables of the tensor operators and their matrix elements. *J. Phys.: Condens. Matter* **16**, 5825-5847.
- Rudowicz C. & Gnutek P. (2009). Modeling techniques for analysis and interpretation of electron magnetic resonance (EMR) data for transition ions at low symmetry sites in crystals - a primer for experimentalists. *Physica B*, **404**, 3582-3582.
- Rudowicz C. & Misra S. K. (2001). Spin-Hamiltonian Formalisms in Electron Magnetic Resonance (EMR) & Related Spectroscopies. *Applied Spectroscopy Reviews* **36**, 11-63.
- Rudowicz C. (1985). Transformation relations for the conventional Okq and normalized O'kq Stevens operator equivalents with with k = 1 to 6 and -k < q < +k. *J. Phys. C: Solid State Phys.*, **18**, 1415-1430; Erratum, ibidem (1985) **18**, 3837.
- Rudowicz C. (1987). Concept of spin Hamiltonian, forms of zero-field splitting and electronic Zeeman Hamiltonians and relations between parameters used in EPR. *A critical review. Magn. Res. Rev.*, **13**, 1-89; Erratum, ibidem (1988) **13**, 335.
- Rudowicz C. (1987). On the derivation of the superposition model formulas using the transformation relations for the Stevens operators. *J. Phys.* **C20**, 6033-7.
- Rudowicz C. (2000). On the relations between the zero-field splitting parameters in the extended Stevens operator notation and the conventional ones used in EMR for orthorhombic and lower symmetry. *J. Phys.: Condensed Matter* **12**, L417-L423.
- Siegel E. & Müller K. A. (1979). Local position of Fe<sup>3+</sup> in ferroelectric BaTiO<sub>3</sub>. *Phys. Rev. B* **20**, 3587-3596.
- Stefaniuk I., Rudowicz C., Gnutek P. & Suchocki A. (2009). EPR Study of Cr<sup>3+</sup> and Fe<sup>3+</sup> Impurity Ions in Nominally Pure and Co<sup>2+</sup>-Doped YAlO<sub>3</sub> Single Crystals. *Appl. Magn. Res.* **36**, 371-380.
- Vasylechko L., Matkovskii A., Savytski D., Suchocki A. & Wallrafen F. (1999). Crystal structure of GdFeO<sub>3</sub>-type rare earth gallates and aluminates. *J. Alloys Compds.* **291**, 57-65.
- Yamaga M., Takeuchi H., Han T. J. & Henderson B. (1993). Electron paramagnetic resonance and optical spectra of Cr<sup>3+</sup>-doped YAlO<sub>3</sub>. *J. Phys.: Condens. Matter* **5**, 8097-8104.
- Yamaga M., Yosida T., Henderson B., O'Donnell K. & Date M. (1992). Electron paramagnetic resonance and optical spectra of Ti<sup>3+</sup>-doped YAlO<sub>3</sub>. *J. Phys.: Condens. Matter* **4**, 7285-7294.
- Yeom T. H., Chang Y. M., Choh S. H. & Rudowicz C. (1994). Experimental and theoretical investigation of spin Hamiltonian parameters for low symmetry Fe<sup>3+</sup> centre in LiNbO<sub>3</sub>. *Physica Status Solidi (b)* **185**, 409-415.